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High precision determination of α_s from a global fit of jet rates

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ABSTRACT: We present state-of-the-art extractions of the strong coupling based on $N^3\text{LO}+\text{NNLL}$ accurate predictions for the two-jet rate in the Durham clustering algorithm at e^+e^- collisions, as well as a simultaneous fit of the two- and three-jet rates taking into account correlations between the two observables. The fits are performed on a large range of data sets collected at the LEP and PETRA colliders, with energies spanning from 35 GeV to 207 GeV. Owing to the high accuracy of the predictions used, the perturbative uncertainty is considerably smaller than that due to hadronization. Our best determination at the Z mass is $\alpha_s(M_Z) = 0.11881 \pm 0.00063(\text{exp.}) \pm 0.00101(\text{hadr.}) \pm 0.00045(\text{ren.}) \pm 0.00034(\text{res.})$, which is in agreement with the latest world average and has a comparable total uncertainty.

KEYWORDS: Jets, QCD Phenomenology

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1 Introduction

Measurements using hadronic final states in e^+e^- annihilation provide a unique opportunity to study the strong interaction in a well controlled environment without strongly interacting particles in the initial state. Accordingly, Quantum Chromodynamics (QCD) was tested extensively at LEP (see e.g. [1–4]).

QCD is a well-established theory by now, and its coupling constant α_s has been measured in a variety of different processes at different energies. Still, the last PDG average of α_s has an uncertainty of about 1% [5], which is considerably larger than errors in other

gauge couplings. This uncertainty has an important impact on the current LHC precision physics program. Furthermore, a number of outlier fits of α_s exist, hence any further precise determination of the value of the strong coupling constant is very valuable.

Many measurements of the strong coupling α_s at e^+e^- colliders are based on comparisons of differential distributions of event shapes or jet rates to perturbative predictions. Presently progress in such measurements depends solely on the improvement in the accuracy of theoretical predictions as new data are not foreseen in the near future.

Compared to event shapes, jet-rates are known to be less sensitive to hadronization corrections,¹ hence more suited for precise determinations of the strong coupling constant. Fully differential predictions for the process $e^+e^- \rightarrow 3$ partonic jets are available to NNLO (α_s^3) accuracy [7–11]. Using the predictions for the three-jet rate at NNLO and the total cross-section at N³LO [12], the two-jet rate can be deduced at N³LO accuracy.

For a jet rate with a resolution parameter y , the fixed-order predictions have a limited range of validity and are not reliable near the boundaries of the phase space, dominated by soft and collinear QCD radiation. In particular, for $y \rightarrow 0$, the perturbative prediction at order $\mathcal{O}(\alpha_s^n)$ features logarithmic terms of the type $\alpha_s^n L^m$ with $L = \ln(y)$, and $m \leq 2n$. When L is large, such terms have to be summed up to all orders in perturbation theory.

The accuracy of the resummation for observables for which double logarithmic terms $\mathcal{O}(\alpha_s^n L^{2n})$ exponentiate, is usually defined in terms of the logarithm of the cumulative distribution Σ . For such exponentiating observables, we define leading logarithms (LL) as terms of the form $\alpha_s^n L^{n+1}$, next-to-leading logarithms (NLL) as $\alpha_s^n L^n$, next-to-next-to-leading logarithms (NNLL) as $\alpha_s^n L^{n-1}$, in $\ln \Sigma$. The state of the art for most event shapes and jet rates measured at LEP is either NNLL or even N³LL [13–24]. These resummed predictions matched to fixed-order ones were used for precise extractions of α_s [13, 14, 17, 25, 26].

Since the structure of logarithmic terms of predictions for jet rates is commonly more involved than that of most event shapes, until very recently, only next-to-double logarithmic corrections $\alpha_s^n L^{2n-1}$ to these observables were known [27]. In this study, we focus on jet rates obtained with the Durham clustering algorithm [27] and will use for the first time NNLL predictions for the two-jet rate $R_2(y)$, which became available in ref. [20].

The main result of this paper is an extraction of the strong coupling from the Durham two-jet rate to data measured at the LEP and PETRA colliders which relies on N³LO+NNLL accurate theoretical predictions. As an additional result, we also present a fit based simultaneously on the two- and three-jet rates, where the latter is computed at NNLO accuracy. For the first time, a fit of the coupling based on the two-jet rate features perturbative uncertainties that are considerably smaller than those of hadronization. This is due to the accurate N³LO+NNLL prediction adopted in the extraction.

2 Observables and predictions

2.1 Jet rates

A jet clustering algorithm is a procedure to classify final-state events into different jet multiplicities. This categorisation depends on the underlying algorithm used. In this

¹See, for instance, the discussion in section 4.5 of ref. [6].

paper we adopt the Durham clustering [27]. This is a sequential recombination algorithm, which requires a distance measure in phase space between momenta² p_i^μ and p_j^μ ,

$$y_{ij} = 2 \frac{\min(E_i^2, E_j^2)}{E_{\text{vis}}^2} (1 - \cos \theta_{ij}), \quad (2.1)$$

where θ_{ij} is the angle between the spatial components of the pair and E_{vis} the visible energy in the event. If the smallest of these, $y_{\min} = \min(\{y_{ij}\})$ is below a pre-defined number, y_{cut} , then the corresponding pair of momenta is recombined into a single one. The procedure continues until all distance measures become larger than y_{cut} . The momenta are recombined using some recombination scheme. Here we adopt the E -scheme [27], according to which the four-momenta of the two clustering particles are simply added together. The n -jet rate is then defined as

$$R_n(y) = \frac{\sigma_{n\text{-jet}}(y)}{\sigma_{\text{tot}}}, \quad (2.2)$$

where $\sigma_{n\text{-jet}}$ is the cross section for n -jet production in hadronic final states obtained with the above algorithm and σ_{tot} is the total hadronic cross section.

In the following sections we briefly review the fixed-order and resummed predictions used in this study. The two results can be combined by means of a matching procedure to obtain the predictions that we ultimately use in the fit.

2.2 Fixed-order predictions

In this work we use predictions obtained with the CoLoRFulNNLO method [11, 28, 29]. The perturbative expansion of the n -jet rate R_n as function of y_{cut} at the default renormalization scale $\mu_{\text{ren}} = Q$ reads

$$R_n(y_{\text{cut}}) = \delta_{2,n} + \frac{\alpha_S(Q)}{2\pi} A_n(y_{\text{cut}}) + \left(\frac{\alpha_S(Q)}{2\pi}\right)^2 B_n(y_{\text{cut}}) + \left(\frac{\alpha_S(Q)}{2\pi}\right)^3 C_n(y_{\text{cut}}) + \mathcal{O}(\alpha_s^4), \quad (2.3)$$

where A , B and C are the perturbative coefficients computed by the MCCSM code [30]. For massless quarks, these coefficients are independent of Q . The renormalization scale dependence of the fixed-order prediction can be restored using the renormalization group equation for α_s ,

$$R_n(y_{\text{cut}}, \mu_{\text{ren}}) = \delta_{2,n} + \frac{\alpha_S(\mu_{\text{ren}})}{2\pi} A_n(y_{\text{cut}}, x_R) + \left(\frac{\alpha_S(\mu_{\text{ren}})}{2\pi}\right)^2 B_n(y_{\text{cut}}, x_R) + \left(\frac{\alpha_S(\mu_{\text{ren}})}{2\pi}\right)^3 C_n(y_{\text{cut}}, x_R) + \mathcal{O}(\alpha_s^4), \quad (2.4)$$

where

$$\begin{aligned} A_n(y_{\text{cut}}, x_R) &= A_n(y_{\text{cut}}), \\ B_n(y_{\text{cut}}, x_R) &= B_n(y_{\text{cut}}) + A_n(y_{\text{cut}}) \frac{1}{2} \beta_0 \ln(x_R^2), \\ C_n(y_{\text{cut}}, x_R) &= C_n(y_{\text{cut}}) + B_n(y_{\text{cut}}) \beta_0 \ln(x_R^2) + A_n(y_{\text{cut}}) \left(\frac{1}{4} \beta_1 \ln(x_R^2) + \frac{1}{4} \beta_0^2 \ln^2(x_R^2) \right), \end{aligned} \quad (2.5)$$

²These momenta belong to either particles, or pseudo-jets obtained during the recombination process.

with $x_R = \mu_{\text{ren}}/Q$, $\beta_0 = (11C_A - 4n_f T_R)/3$ and $\beta_1 = (34C_A^2 - 20C_A T_R n_f - 12C_F T_R n_f)/3$. The numerical values for fixed-order coefficients for R_3 , R_4 and R_5 are reported in appendix A, table 3. These can be used to build up the necessary fixed-order predictions for R_2 that we use in the following.

2.3 Resummed predictions

The resummation technique adopted here was formulated in refs. [19, 20, 23], hence we present only the main features of these results. The essence of the procedure described in ref. [19] is that the NLL cross section is given by all-order configurations made of partons independently emitted off the Born legs and widely separated in angle [31]. The NNLL corrections are obtained by correcting a *single* parton of the above ensemble to account for all kinematic configurations that give rise to NNLL effects [20]. The two-jet rate at NNLL can be written as

$$R_2(y_{\text{cut}}) = e^{-R_{\text{NNLL}}(y_{\text{cut}})} \left[\left(1 + \frac{\alpha_s(Q)}{2\pi} H^{(1)} + \frac{\alpha_s(Q\sqrt{y_{\text{cut}}})}{2\pi} C_{\text{hc}}^{(1)} \right) \mathcal{F}_{\text{NLL}}(y_{\text{cut}}) + \frac{\alpha_s(Q)}{\pi} \delta\mathcal{F}_{\text{NNLL}}(y_{\text{cut}}) \right], \quad (2.6)$$

where the Sudakov radiator R_{NNLL} and the coefficients $H^{(1)}$, $C_{\text{hc}}^{(1)}$ are defined in ref. [23],³ while the functions \mathcal{F}_{NLL} and $\delta\mathcal{F}_{\text{NNLL}}$ are given in refs. [19, 20].

For the two-jet rate R_2 the resummation is performed with the **ARES** program [20] and the matching to fixed-order is done according to the $\ln R$ scheme [32].

The resummation of the three-jet rate is much more involved due to the extra number of emitting particles. Accordingly, the state-of-the-art resummed predictions have a much lower logarithmic accuracy and include only terms $\mathcal{O}(\alpha_s^n L^{2n})$ and $\mathcal{O}(\alpha_s^n L^{2n-1})$ in $R_3(y)$ [27], in contrast to the logarithmic counting that we introduced for R_2 which refers to the logarithm of the jet rate. While R_3 is more sensitive to α_s than R_2 , the low theoretical accuracy of the resummation does not guarantee a good theoretical control in the region where the logarithms are large. Therefore, for the present analysis, we do not perform any resummation of R_3 and limit the fit to a range where the fixed-order is reliable.

2.4 Effects of quark masses

The effect of the non-vanishing b -quark mass on the predictions has also been considered in the literature. In particular, predictions for $e^+e^- \rightarrow$ partons are known including α_s^2 corrections with massive b -quarks [33]. The resummed predictions for the Durham R_2 and R_3 observables with non-zero b -quark masses are only known at next-to-double logarithmic accuracy [34] (i.e. $\alpha_s^n L^{2n-1}$ in the jet-rates) and are not used in this analysis as this does not match our target accuracy needed to guarantee a robust theoretical control even in the region where the logarithms become large.

In order to take b -quark mass corrections into account, we subtract the fraction of b -quark events $r_b(Q)$ from the massless result and add back the corresponding massive

³This observable corresponds to setting $a = 2$ and $b_\ell = 0$ in the corresponding equations.

contribution computed at fixed order. Hence, we include mass effects directly at the level of the final distributions according to the formulae

$$\begin{aligned} R_2(y) &= R_2^{\text{N}^3\text{LO}+\text{NNLL}}(y)_{m_b=0}(1-r_b(Q)) + r_b(Q)R_2^{\text{NNLO}}(y)_{m_b \neq 0}, \\ R_3(y) &= R_3^{\text{NNLO}}(y)_{m_b=0}(1-r_b(Q)) + r_b(Q)R_3^{\text{NLO}}(y)_{m_b \neq 0}. \end{aligned} \quad (2.7)$$

The latter quantities are obtained by combining the total cross section at NNLO including mass corrections as obtained from ref. [35], and the three- and four-jet rate $\mathcal{O}(\alpha_s^2)$ predictions as computed with the Zbb4 program [33, 36]. The strong coupling used in the prediction of Zbb4 is then converted into a 5-flavour coupling by means of the matching relation for $\alpha_s(m_b)$ [37, 38].

We define the fraction of b -quark events as the ratio of the total b -quark production cross section divided by the total hadronic cross section,

$$r_b(Q) \equiv \frac{\sigma_{m_b \neq 0}(e^+e^- \rightarrow b\bar{b})}{\sigma_{m_b \neq 0}(e^+e^- \rightarrow \text{hadrons})}. \quad (2.8)$$

We evaluate the ratio of these cross sections to approximate $\mathcal{O}(\alpha_s^3)$ according to ref. [35].

For the bottom quark we used a pole mass of $m_b = 4.78 \text{ GeV}$, which is consistent with the corresponding world average [39].

3 Determination of α_s

To extract the strong coupling we compare the theory predictions described above to the available data, taking into account the non-perturbative corrections from Monte Carlo (MC) models, as described in detail in the following.

3.1 Data sets

We select experimental data that satisfy the following basic requirements: (i) measurements are obtained with both charged and neutral final state particles, (ii) corrections for detector effects have been taken into account, and (iii) corrections for initial-state QED radiation have been taken into account. We found that the data from refs. [40–48] satisfy these basic criteria.

However, the measurements of ref. [41] and ref. [42] are superseded by the measurements in ref. [43] and hence not included in our analysis. The analyses in ref. [44] was excluded as the provided combined uncertainties are much lower than those from later refined analyses and are close to expected statistical uncertainties estimated from earlier analysis in ref. [49]. We also excluded the measurements from ref. [45] as these contain explicit corrections for the contributions of the process $e^+e^- \rightarrow b\bar{b}$ from MC simulations and cannot be treated in the same way as the data without such corrections.

From the selected analyses we also excluded the measurements at $\sqrt{s} = 172 \text{ GeV}$ from ref. [46] as the background subtraction procedure performed with data of limited statistics has biased the measurement, e.g. introduced a non-monotonous behaviour of $R_2(y)$, see ref. [50] for details. For similar reasons we exclude data from ref. [47]. The sum of the rates

Experiment	Data \sqrt{s} , GeV	MC \sqrt{s} , GeV	Events
OPAL [43]	91.2(91.2)	91.2	1508031
OPAL [43]	189.0(189.0)	189	3300
OPAL [43]	183.0(183.0)	183	1082
OPAL [43]	172.0(172.0)	172	224
OPAL [43]	161.0(161.0)	161	281
OPAL [43]	130.0–136.0(133.0)	133	630
L3 [46]	201.5–209.1(206.2)	206	4146
L3 [46]	199.2–203.8(200.2)	200	2456
L3 [46]	191.4–196.0(194.4)	194	2403
L3 [46]	188.4–189.9(188.6)	189	4479
L3 [46]	180.8–184.2(182.8)	183	1500
L3 [46]	161.2–164.7(161.3)	161	424
L3 [46]	135.9–140.1(136.1)	136	414
L3 [46]	129.9–130.4(130.1)	130	556
JADE [43]	43.4–44.3(43.7)	44	4110
JADE [43]	34.5–35.5(34.9)	35	29514
ALEPH [48]	91.2(91.2)	91.2	3600000
ALEPH [48]	206.0(206.0)	206	3578
ALEPH [48]	189.0(189.0)	189	3578
ALEPH [48]	183.0(183.0)	183	1319
ALEPH [48]	172.0(172.0)	172	257
ALEPH [48]	161.0(161.0)	161	319
ALEPH [48]	133.0(133.0)	133	806

Table 1. Data used for the determination of α_s in this work. The ranges of collision energies, their weighted average value (in brackets) and the number of events for each experiment are given as quoted in the original publications.

for the data sets in ref. [48] deviates from unity, and the largest deviation reaches almost 0.03 for the data set at $\sqrt{s} = 200$ GeV. For this reason we excluded this data from the fits as well.

We summarise the information on the selected data sets in table 1. For some data sets the uncertainties were updated before the fit, as follows. We added in quadrature the two available systematic uncertainties for the measurements at $\sqrt{s} = 91$ GeV from ref. [48]. The data from JADE [43] does not include systematic uncertainties related to the choice of MC samples used for the calculations of detector corrections. Later studies with the same data [45] indicated that such uncertainties can be at least as large as the statistical uncertainties. Therefore, in this case we added a relative uncertainty of 1.5% as an estimation of missing systematic uncertainties.

The measurements of the jet rates selected for the analysis are provided in the original publications without correlations between the individual points and without decomposition of the total systematic uncertainties. To perform an accurate extraction procedure, we examined the available data and uncertainties and built a covariance matrix for all measured sets of data. This procedure consists of the following steps.

In the first step we estimated the statistical correlation matrix of the individual points in the fit range as described in ref. [51] from the MC generated samples. We find that these are in good agreement with statistical correlation matrices obtained from data in unpublished ref. [52]. In the second step we built the full covariance matrix for each measurement from the correlation matrix, statistical and systematic uncertainties.

As the original publications do not contain enough information, the procedure is based on some assumptions. Namely, for a pair of measurements $R_n(y_1)$ and $R_m(y_2)$ we make the following assumptions on the correlation coefficients of the systematic uncertainties: (i) $\text{corr}_{\text{syst}}[R_n(y_1), R_m(y_2)] = \text{corr}_{\text{stat}}[R_n(y_1), R_m(y_2)] \times K \times \rho^{|\log(y_1) - \log(y_2)|}$ for $|\log(y_1) - \log(y_2)| < 0.25$ and (ii) $\text{corr}_{\text{syst}}(R_n(y_1), R_m(y_2)) = 0$ for $|\log(y_1) - \log(y_2)| > 0.25$. We selected $\rho = 0.006$ in order to mimic patterns of systematic uncertainties observed in ref. [52]. This corresponds to $\text{corr}_{\text{syst}}[R_n(y_1), R_m(y_2)] \approx 0.5K \text{corr}_{\text{stat}}[R_n(y_1), R_m(y_2)]$ for measurements with $|\log(y_1) - \log(y_2)| = 0.125$, i.e. in the neighbouring bins for typical binning used in the measurements. K is equal to 1 if $n = m$ and 0.5 otherwise. This approach approximates correlations between observables R_2 and R_3 .

3.2 Monte Carlo event generation setup

In our analysis we model non-perturbative effects in the $e^+e^- \rightarrow \text{hadrons}$ process using state-of-the-art particle level MC event generators. As usual, we estimate the non-perturbative corrections of the jet rate distributions by comparing distributions at hadron and parton level in the simulated samples. In particular, we used the **Herwig7.1.4** [53] MC event generator to deliver our final results, and the **Sherpa2.2.6** [54, 55] MC event generator for cross-checks.

We generated event samples for the process $e^+e^- \rightarrow \text{hadrons}$ at the centre-of-mass energies listed in table 1. In all cases, we switched off the simulation of initial state radiation and used default generator settings, unless stated otherwise. We used $\alpha_s(M_Z) = 0.1181$ [39] as input for the generation. Moreover, we adopted the G_μ scheme with input parameters $M_Z = 91.1876 \text{ GeV}$, $M_W = 80.379 \text{ GeV}$ and $G_F = 1.1663787 \times 10^{-5} \text{ GeV}^{-2}$. The pole masses of b - and t -quarks were set to 4.78 GeV and 173 GeV respectively.

The **Herwig7.1.4** samples were generated with the unitarised MENLOPS method [56] using the **MadGraph5** [57] matrix element generator and the **OpenLoops** [58] one-loop library to produce matrix elements for the 2, 3, 4 and 5 parton final states in the hard process. The two- and three-parton final-state matrix elements have again NLO accuracy in perturbative QCD and the QCD matrix elements were calculated using massive b -quarks. The merging parameter was set to $\sqrt{s} \times 10^{-1.25}$.

To test the fragmentation and hadronization model dependence, the events generated with **Herwig7.1.4** were hadronized using either the cluster fragmentation model [59] or the Lund string fragmentation model [60]. The cluster fragmentation model is natively implemented in **Herwig7.1.4**. The Lund string fragmentation model is implemented in **Pythia8.2.35** [61] and was used in **Herwig7.1.4** via the **ThePEG1.6.1++** [62] toolkit. For both setups the hadron decays were performed by the **EvtGen1.7.0** [63] package. We label the first setup H^C and the second one H^L .

The Sherpa2.2.6 samples were generated with the MENLOPS method using the matrix element generators AMEGIC [64], COMIX [65] and the OpenLoops [58] one-loop library to produce matrix elements of the processes $e^+e^- \rightarrow Z/\gamma^* \rightarrow 2, 3, 4, 5$ partons. The two-parton final state-matrix elements again have NLO QCD accuracy. The QCD matrix elements were calculated assuming massive b -quarks. The merging parameter Y_{cut} was set to $10^{-2.5}$. The events were hadronized using the cluster fragmentation model [66] as implemented in Sherpa2.2.6. We label this setup S^C .

3.3 Estimation of hadronization effects from MC models

The estimation of hadronization corrections is an integral part of comparing the parton-level QCD predictions to the data measured at hadron (particle) level. While the principle of local parton-hadron duality leads to close values of observable quantities at parton and hadron level, the difference between them is not negligible and must be taken into account if one aims at a precise determination of α_s . One possible way to do this is to apply correction factors estimated from MC simulations to the perturbative predictions.

To obtain the jet rates at parton and hadron level, we processed the MC generated samples in the same way as for the data, using undecayed/stable particles for hadron level calculations. To define jets with the Durham clustering, we used the implementation of the FastJet/fjcore3.3.0 package [67]. The selected resulting distributions are shown in figure 8 in appendix B.

The predictions obtained with all setups describe the data well for all ranges of y , with the exception of the regions of small y at all values of \sqrt{s} (see figures 9, 10 and 11).

Among all considered MC setups, H^L was selected to be the reference. The selected setup uses the very well tested Lund hadronization model that provides stable and physically reliable predictions throughout a wide range of centre-of-mass energies in e^+e^- collisions. Moreover, precisely this hadronization model was used for the modelling of e^+e^- collisions in the original publications [43, 46, 48].

To estimate the uncertainty related to the hadronization modelling, we consider half the difference between the H^L and H^C setups, as explained in more detail in the following.

MC generators were already used to model hadronization effects in previous QCD analyses of e^+e^- data [26, 68]. Typically, hadron level predictions for every observable were obtained by multiplying the perturbative predictions by some factor derived from the analysis of MC-generated events. In the present analysis the approach was amended to take into account physical constraints on R_n , namely that jet rates are positive and that their sum should be one. These constraints are implemented by introducing the variables ξ_1 and ξ_2 such that at parton level

$$R_2^{(p)} = \cos^2 \xi_1, \quad R_3^{(p)} = \sin^2 \xi_1 \cos^2 \xi_2 \quad \text{and} \quad R_{\geq 4}^{(p)} = \sin^2 \xi_1 \sin^2 \xi_2, \quad (3.1)$$

with the constraint

$$R_2^{(p)} + R_3^{(p)} + R_{\geq 4}^{(p)} = 1. \quad (3.2)$$

The corresponding relations at hadron level are

$$\begin{aligned} R_2^{(h)} &= \cos^2(\xi_1 + \delta\xi_1), \quad R_3^{(h)} = \sin^2(\xi_1 + \delta\xi_1) \cos^2(\xi_2 + \delta\xi_2), \\ R_{\geq 4}^{(h)} &= \sin^2(\xi_1 + \delta\xi_1) \sin^2(\xi_2 + \delta\xi_2). \end{aligned} \quad (3.3)$$

The functions $\delta\xi_1(y)$ and $\delta\xi_2(y)$ take into account the non-perturbative corrections. We obtained the numerical values for $\delta\xi_1(y)$ and $\delta\xi_2(y)$ from the MC simulated samples. In order to extract the hadronization corrections we proceed as follows. For a given y bin, we extract ξ_1 from the parton level prediction for the two-jet rate, and ξ_2 from the same prediction for the three-jet rate. The shifts $\delta\xi_1$ and $\delta\xi_2$ are then extracted in the same fashion from the hadron level predictions. The extracted values and corresponding interpolated functions (splines) are shown in figure 12 in appendix C for selected centre-of-mass energies. The size of the derived hadronization corrections can be read off figure 13. As expected, we see that hadronization corrections increase at small values of the two- and three-jet rates and that the corrections become less important at higher energies.

4 Fit procedure

To find the optimal value of α_s , we used the MINUIT2 [69] package to minimize $\chi^2 = \sum_{\text{data sets}} \chi^2(\alpha_s)_{\text{data set}}$, where $\chi^2(\alpha_s)$ was calculated for each data set as

$$\chi^2(\alpha_s) = \vec{r} V^{-1} \vec{r}^T, \quad \vec{r} \equiv (\vec{D} - \vec{P}(\alpha_s)), \quad (4.1)$$

where \vec{D} stands for the vector of data points, $\vec{P}(\alpha_s)$ is the vector of theoretical predictions, and V is the covariance matrix for the experimental data \vec{D} .

We choose the fit range as follows. In order to assure that the implementation of the hadronization corrections is unitary, i.e. the constraint of eq. (3.2) is satisfied, we set the upper bound of the fit range below the kinematical limit for four jet production, $\log_{10}(y) = \log_{10}(1/6) \simeq -0.8$. We therefore choose $\log_{10}(y) = -1$ as an upper bound. Moreover, we adapt the lower bound to the centre-of-mass energy in order to take into account that hadronization corrections become more important at lower energies. Accordingly, we fix the lower bound $\log_{10} y_{\min}(Q)$ of the fit range as $\log_{10} y_{\min}(Q) = \log_{10} y_{\min}(M_Z) + \mathcal{L}$ with $\mathcal{L} = \log_{10}(M_Z^2/Q^2)$. Different values used for $\log_{10}(y_{\min}(M_Z))$ for R_2 and R_3 are indicated in the first columns in tables 2, 4 and 5.

4.1 Fit of the coupling with the two-jet rate R_2

To obtain the most precise results, we first concentrate on fits that include the two-jet rate R_2 solely. The results of the fits, together with the used fit ranges, are given in table 2. We show the results obtained via a fit of R_2 both at N³LO and N³LO+NNLL, with different hadronization models. The corresponding values of χ^2 divided by the number of degrees of freedom (ndof) in the global fit is also reported. The values χ^2/ndof in the global fit as well as the values for every data set (not shown) are all of order unity, which can be viewed as a support for our correlation model. These values can be compared to the ones obtained in similar analyses. For instance, in ref. [70], the χ^2/ndof for fits with statistical uncertainties only varies for different observables between 0.5 and 60.

Fit ranges, $\log_{10}(y)$ Hadronization	N ³ LO $\chi^2/ndof$	N ³ LO+NNLL $\chi^2/ndof$
$[-1.75 + \mathcal{L}, -1]$ S^C	0.12121 ± 0.00095 20/86 = 0.24	0.11849 ± 0.00092 20/86 = 0.24
$[-2 + \mathcal{L}, -1]$ S^C	0.12114 ± 0.00081 26/100 = 0.26	0.11864 ± 0.00075 26/100 = 0.26
$[-2.25 + \mathcal{L}, -1]$ S^C	0.12119 ± 0.00060 44/150 = 0.29	0.11916 ± 0.00063 44/150 = 0.29
$[-2.5 + \mathcal{L}, -1]$ S^C	0.12217 ± 0.00052 89/180 = 0.50	0.12075 ± 0.00055 107/180 = 0.59
$[-1.75 + \mathcal{L}, -1]$ H^C	0.11957 ± 0.00098 22/86 = 0.26	0.11698 ± 0.00093 22/86 = 0.25
$[-2 + \mathcal{L}, -1]$ H^C	0.11923 ± 0.00079 29/100 = 0.29	0.11687 ± 0.00076 28/100 = 0.28
$[-2.25 + \mathcal{L}, -1]$ H^C	0.11868 ± 0.00068 43/150 = 0.28	0.11679 ± 0.00064 40/150 = 0.27
$[-2.5 + \mathcal{L}, -1]$ H^C	0.11849 ± 0.00050 58/180 = 0.32	0.11723 ± 0.00053 58/180 = 0.32
$[-1.75 + \mathcal{L}, -1]$ H^L	0.12171 ± 0.00109 21/86 = 0.25	0.11897 ± 0.00092 21/86 = 0.24
$[-2 + \mathcal{L}, -1]$ H^L	0.12144 ± 0.00078 28/100 = 0.28	0.11893 ± 0.00075 26/100 = 0.26
$[-2.25 + \mathcal{L}, -1]$ H^L	0.12080 ± 0.00069 43/150 = 0.28	0.11881 ± 0.00063 39/150 = 0.26
$[-2.5 + \mathcal{L}, -1]$ H^L	0.12024 ± 0.00051 57/180 = 0.32	0.11897 ± 0.00053 52/180 = 0.29

Table 2. Fit of $\alpha_s(M_Z)$ from experimental data for R_2 obtained using N³LO and N³LO+NNLL predictions, three different hadronization models and four different choices of the fit range, as given in the brackets, with $\mathcal{L} = \log_{10}(M_Z^2/Q^2)$. The reported uncertainty is the fit uncertainty as given by MINUIT2.

From the results given in the table one notices that the effect of the resummation is to move the fitted $\alpha_s(M_Z)$ to slightly lower values. The quality of the fit, with or without resummation, is very similar. The benefit of including the resummation will become evident when perturbative uncertainties of the results are discussed.

As our best fit we quote the result obtained from the fits of the R_2 observable with the H^L hadronization model in the fit range $[-2.25 + \mathcal{L}, -1]$, that reads

$$\alpha_s(M_Z) = 0.11881 \pm 0.00063(\text{exp.}) \pm 0.00101(\text{hadr.}) \pm 0.00045(\text{ren.}) \pm 0.00034(\text{res.}), \quad (4.2)$$

where the quoted uncertainties are coming from MINUIT2 (exp.), variation of renormalization scale (ren.), variation of resummation scale (res.) and choice of hadronization model (hadr.). The estimation of these uncertainties is described in the following subsection.

Finally, we show in figures 1, 2 and 3 the comparison of data at different energies with theory predictions using $\alpha_s(M_Z)$ obtained from our global fit, eq. (4.2).

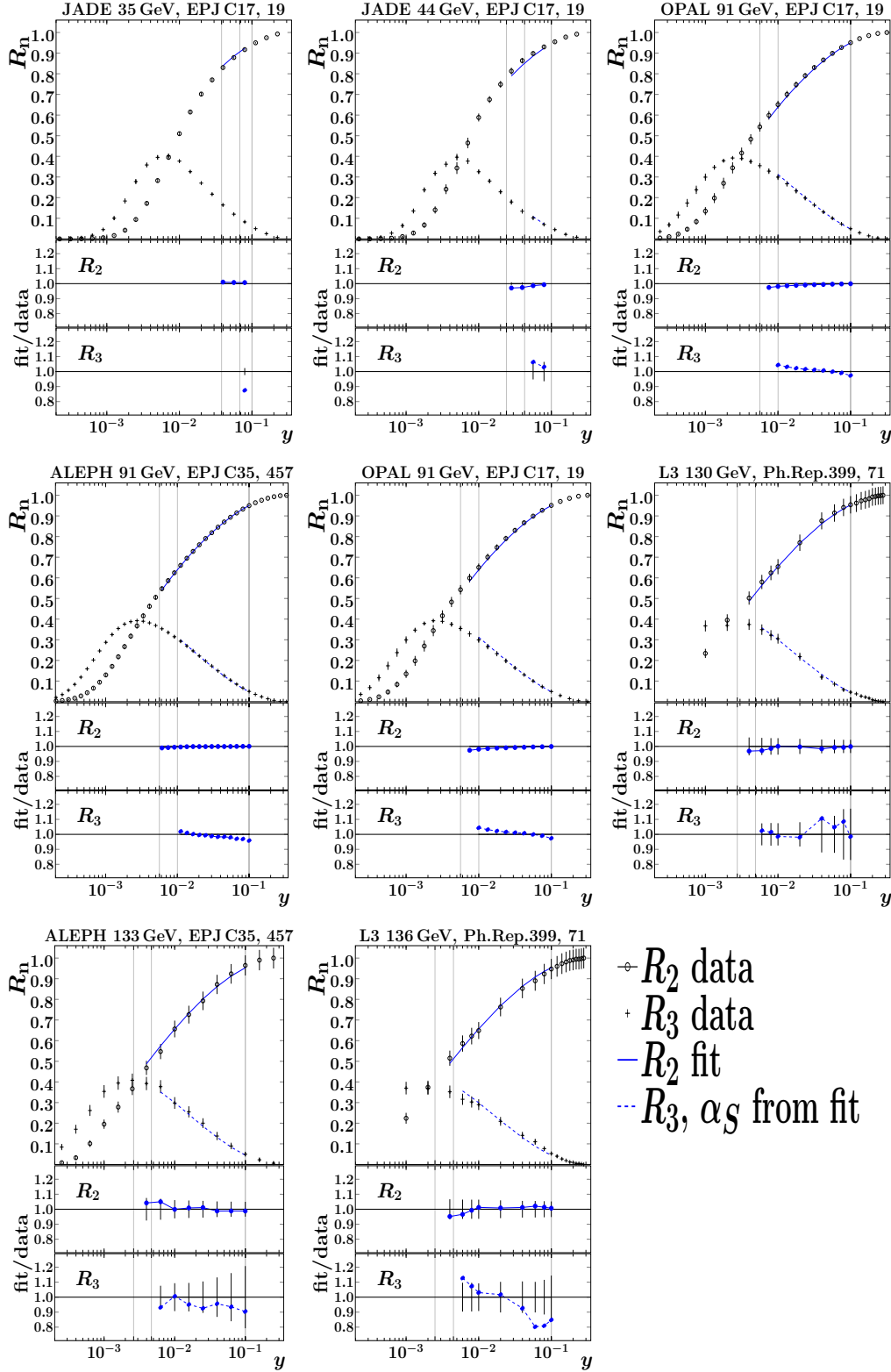


Figure 1. Comparison of data and perturbative predictions supplemented by hadronization corrections in the H^L model using for the strong coupling the value obtained from our global fit, eq. (4.2).

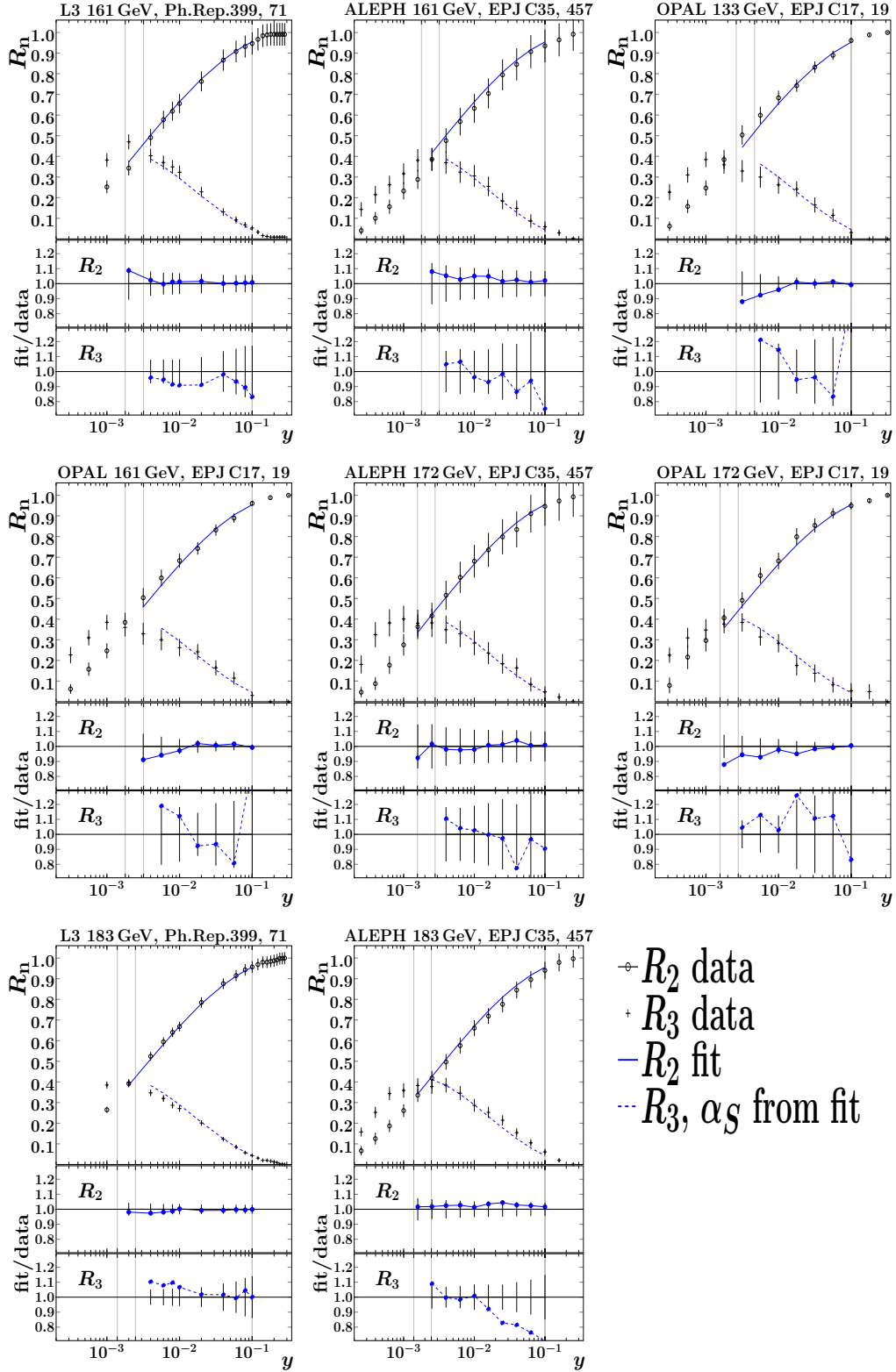


Figure 2. Comparison of data and perturbative predictions supplemented by hadronization corrections in the H^L model using for the strong coupling the value obtained from our global fit, eq. (4.2).

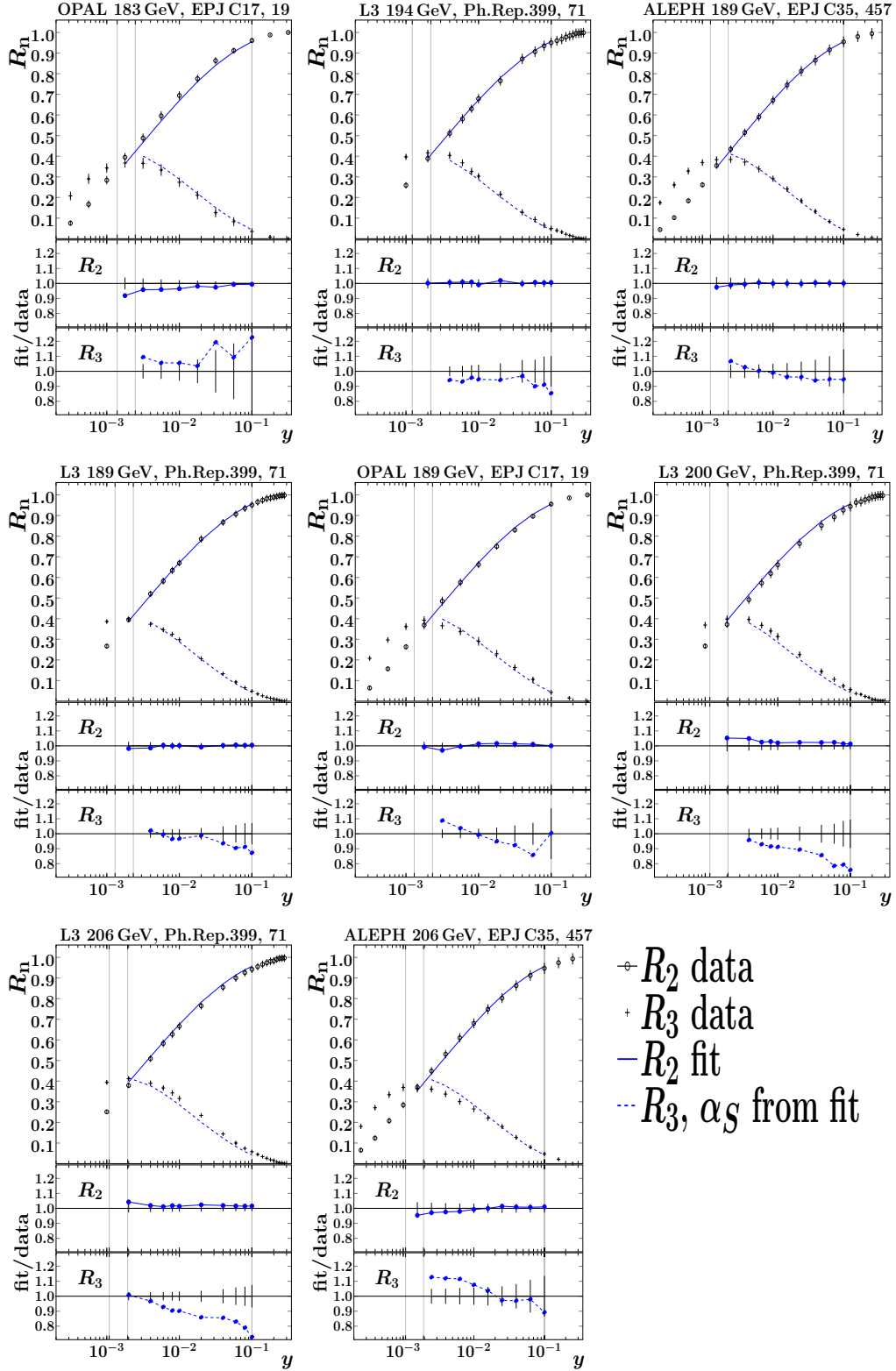


Figure 3. Comparison of data and perturbative predictions supplemented by hadronization corrections in the H^L model using for the strong coupling the value obtained from our global fit, eq. (4.2).

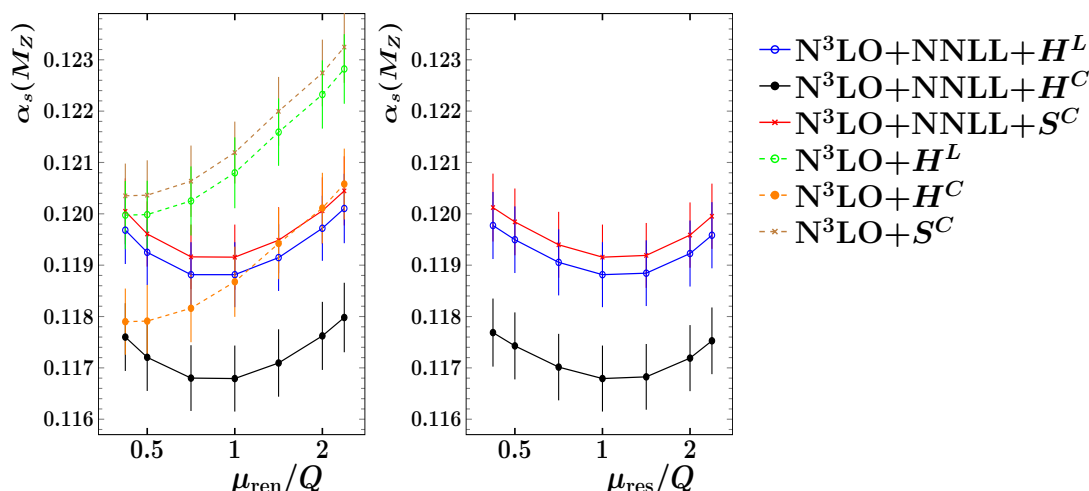


Figure 4. Dependence of (R_2) fit results on the renormalization and resummation scales. The fit range for S^C , H^C and H^L setups is $[-2.25 + \mathcal{L}, -1]$ with $\mathcal{L} = \log_{10}(M_Z^2/Q^2)$.

4.2 Estimation of uncertainties

The systematic uncertainties in α_s are determined following the procedure of [71]. To estimate the size of higher-order terms in the perturbative prediction, we vary the renormalization scale in the range $\mu_{\text{ren}} = Q/2$ and $\mu_{\text{ren}} = 2Q$. Moreover, while keeping $\mu_{\text{ren}} = Q$ fixed, we vary the resummation scale in the range $\mu_{\text{res}} = Q/2$ and $\mu_{\text{res}} = 2Q$. The effects of the individual variations are displayed in figure 4, where different hadronization models are compared. We notice that when resummation is included, a much reduced dependence on the renormalization scale is observed.

The bias due to the selection of the hadronization model is studied by means of the difference between the H^L and H^C setups, see figure 4. In particular, considering the results from the Lund string and cluster hadronization models, the desired systematic uncertainty is obtained as half of the difference between the $\alpha_s(M_Z)$ results obtained in nominal fits with H^L and H^C setups. The obtained numerical value is close to 0.001 (i.e. slightly below 1%). This can be briefly compared to previous estimations obtained with Monte Carlo event generator models in similar analyses. Namely, the values 0.001 [68, 70] and 0.0005 [72] obtained previously allow us to validate our estimation. We stress that we have not performed any tuning of the adopted hadronization models to the data in order to artificially reduce the related uncertainties. This leads to a more conservative, and thus more robust estimate of hadronization uncertainties.

The final uncertainty is obtained by combining each of the above uncertainties in quadrature.

5 Validation of the procedure and further fits

In this section we perform some consistency checks to validate the fitting procedure used above. Moreover, we present an extraction of $\alpha_s(M_Z)$ from a simultaneous fit of R_2 at $N^3\text{LO}+\text{NNLL}$ and R_3 at NNLO.

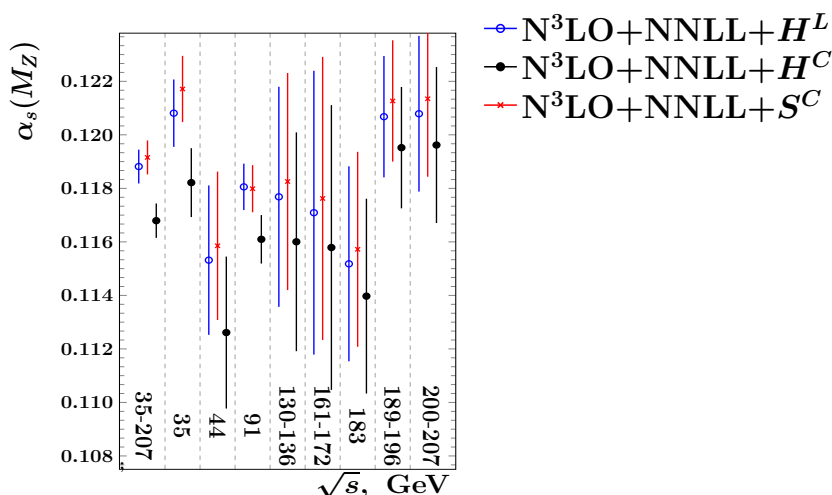


Figure 5. Dependence of (R_2) fit results on the \sqrt{s} of measurement. The fit range is $[-2.25+\mathcal{L}, -1]$. Only statistical uncertainties are shown.

5.1 Fit consistency tests

We have performed a number of consistency tests, as outlined in the following:

1. We repeat the nominal fits in different ranges of \sqrt{s} separately, instead of simultaneously. The results of these fits are shown in figure 5. We do not observe any significant dependence of the results on the centre-of-mass energy.
2. We repeat the nominal fit for R_2 by implementing the hadronization corrections on a bin-by-bin basis as $R_{2,\text{hadron}} = R_{2,\text{parton}} f_2(y)$, where $f_2(y)$ is derived from the MC generated samples. Within this scheme, we find

$$\alpha_s(M_Z) = 0.11881 \pm 0.00063(\text{exp.}) \pm 0.00109(\text{hadr.}).$$

This value is close to the reference result, but the hadronization uncertainty estimated from this setup is more sensitive to changes in the fit range.

3. To test the reliability of the correlation model used for the systematic uncertainties in the reference fit, we use the OPAL data and systematic shifts (uncertainties) from ref. [52]. With this data we perform the fits using the χ^2 definition from eq. (4.1) and the definition that explicitly includes set of N nuisance parameters $B = \{b_1, b_2 \dots b_N\}$ and vectors of systematic shifts (uncertainties) $\vec{S}_1, \vec{S}_2 \dots \vec{S}_N$:⁴

$$\chi^2(\alpha_s, B) = \left(\vec{r} - \sum_{i=1}^N b_i \vec{S}_i \right) V^{-1} \left(\vec{r} - \sum_{i=1}^N b_i \vec{S}_i \right)^T + \sum_{i=1}^N b_i^2. \quad (5.1)$$

⁴Such a treatment of systematic shifts (uncertainties) is widely used in QCD analyses, see ref. [73] as an example.

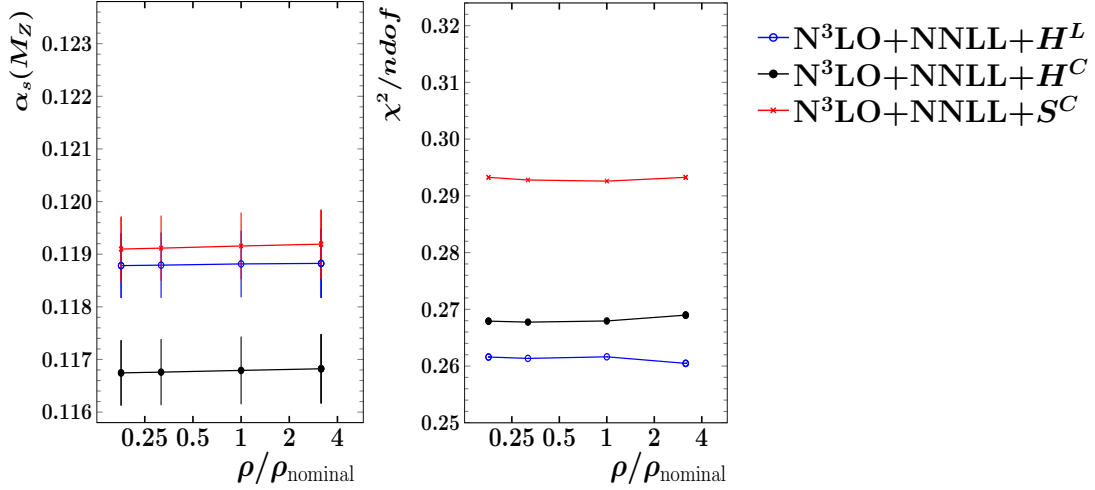


Figure 6. Dependence of the (R_2) fit results on the parameter ρ used in the correlation model. The fit range is fixed to $[-2.25 + \mathcal{L}, -1]$.

The result is

$$\alpha_s(M_Z) = 0.11893 \pm 0.00137(\text{exp.}) \quad \text{for the definition in eq. (4.1)}$$

and

$$\alpha_s(M_Z) = 0.11761 \pm 0.00179(\text{exp.}) \quad \text{for the definition in eq. (5.1).}$$

These results are in fair agreement, and the corresponding values of χ^2 are close, which demonstrates the relatively good performance of the selected correlation model.

4. We repeat the reference fit without the data sets from JADE, for which the explicit verification of correlation model using ref. [52] was not possible. The obtained result $\alpha_s(M_Z) = 0.11838 \pm 0.00077(\text{exp.})$ agrees well with the reference value $\alpha_s(M_Z) = 0.11881 \pm 0.00063(\text{exp.}) \pm 0.00101(\text{hadr.}) \pm 0.00045(\text{ren.}) \pm 0.00034(\text{res.})$.
5. To estimate the error due to the correlation model, we vary the value of the parameter ρ in a wide range. No significant change of fit results is observed, see figure 6.
6. We perform variations of the renormalization and resummation scales simultaneously, as done in ref. [45]. The results are reported in figure 7 and show smaller uncertainty than in the case of independent variations. A similar behaviour was observed in ref. [26].
7. To study the reliability of the theory uncertainty estimation based on the variation of the resummation and renormalization scales, the reference fit using the R_2 observable was repeated using the NNLO+NLL theory predictions, obtaining

$$\alpha_s(M_Z) = 0.11999 \pm 0.00067(\text{exp.}) \pm 0.00105(\text{hadr.}) \pm 0.00210(\text{ren.}) \pm 0.00172(\text{res.}).$$

The corresponding uncertainty band contains the central value of the $N^3\text{LO}+\text{NNLL}$ fit, which demonstrates that the estimated theory error is reliable.

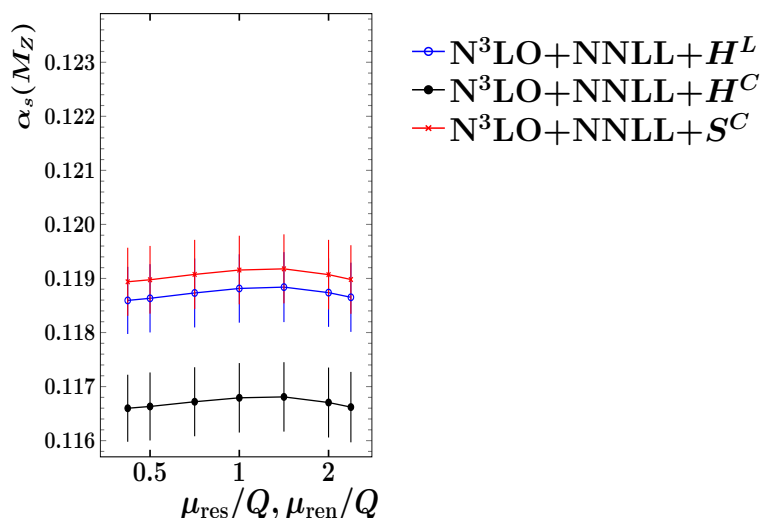


Figure 7. Dependence of the (R_2) fit results on the simultaneous variation of renormalization and resummation scales. The fit range is fixed to $[-2.25 + \mathcal{L}, -1]$.

5.2 Simultaneous fit of the coupling with the two- and three- jet rates R_2 and R_3

An alternative fit was performed with R_2 and R_3 observables simultaneously, with R_3 computed at NNLO. The obtained result using the $[-2.25 + \mathcal{L}, -1][-2 + \mathcal{L}, -1]$ fit ranges for R_2 and R_3 , respectively, and the H^L setup is

$$\alpha_s(M_Z) = 0.11989 \pm 0.00045(\text{exp.}) \pm 0.00098(\text{hadr.}) \pm 0.00046(\text{ren.}) \pm 0.00017(\text{res.}).$$

Unlike the results of the reference fits, the obtained result is sensitive to the selected fit range, see table 4 in appendix D. Taking this effect into account would result in another uncertainty of order 0.001 that is not included in the uncertainties given above.

As a final cross-check we perform a fit for a single point of ALEPH R_3 data at $y = 0.02$ without resummation. The obtained result (with statistical uncertainties only), $\alpha_s(M_Z) = 0.11905 \pm 0.00251(\text{exp.})$, can be compared to the results from ref. [74], $\alpha_s(M_Z) = 0.1175 \pm 0.0020(\text{exp.}) \pm 0.0015(\text{theo.})$. The results agree well despite the differences in the implementation of the hadronization corrections.

6 Summary

The main result of this paper is a first fit of the strong coupling for the two-jet rate that relies on very accurate ($N^3\text{LO}+\text{NNLL}$) predictions. Our main result reads

$$\alpha_s(M_Z) = 0.11881 \pm 0.00063(\text{exp.}) \pm 0.00101(\text{hadr.}) \pm 0.00045(\text{ren.}) \pm 0.00034(\text{res.}).$$

The uncertainty on α_s induced by scale variations is now considerably smaller than that related to hadronization modelling. This is not the case if the fit is performed using only $N^3\text{LO}$ predictions, see figure 4. Furthermore the experimental uncertainty is now comparable to the perturbative one.

After combining the uncertainties in quadrature we obtain

$$\alpha_s(M_Z) = 0.11881 \pm 0.00131(\text{comb.}),$$

where the largest estimation bias comes from the used hadronization model.

Our results agrees with the last PDG average

$$\alpha_s(M_Z)_{\text{PDG2018}} = 0.1181 \pm 0.0011$$

with an uncertainty that is of the same size.

We have also performed a combined fit of the two- and three-jet rates, taking for the first time the correlation between these observables into account. The results of the two fits are fully compatible. However, the fit including R_3 shows a stronger dependence on the fit range than the results of our reference fit based on R_2 only. An accurate resummation for the R_3 observable could potentially reduce the sensitivity to the fit range selection and lead to an even more precise determination of $\alpha_s(M_Z)$.

Acknowledgments

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A Perturbative ingredients

We report in table 3 the numerical results for the fixed-order coefficients introduced in eq. (2.5).

y	A_3	A_4	A_5	B_3	B_4	C_3
0.000249	64.853	2226.9	52659	−2276.3	−18059	−52708
0.000304	61.462	1981.9	43613	−1951.9	−9434	−49312
0.000371	57.969	1745.7	35503	−1644.5	−2392	−45862
0.000453	54.589	1531.6	28680	−1370.8	2875.1	−41808
0.000553	51.315	1337.6	22984	−1126.4	6726.2	−38197
0.000676	48.149	1162.6	18248	−911.38	9293.8	−33891
0.000825	45.081	1005.5	14362	−722.69	10902	−29909
0.001008	42.125	864.92	11185	−558.20	11665	−25860
0.001231	39.267	739.71	8612	−416.39	11905	−22437
0.001503	36.514	628.56	6548	−295.16	11595	−18881
0.001836	33.865	530.39	4910	−192.34	10949	−15681
0.002243	31.318	444.20	3627	−106.70	10084	−12771
0.002739	28.874	368.93	2634	−36.132	9080.6	−10373
0.003346	26.532	303.65	1875	20.7909	8018.6	−8011
0.004087	24.292	247.42	1307	65.2347	6936.3	−6144
0.004992	22.155	199.36	888.5	98.9217	5886.0	−4478
0.006097	20.118	158.62	586.9	123.332	4897.6	−3090
0.007447	18.183	124.43	374.8	139.397	3982.9	−1863
0.009095	16.354	96.058	230.1	148.398	3173.8	−944.6
0.011109	14.621	72.787	134.4	151.554	2469.0	−218.7
0.013569	12.988	53.978	74.07	149.800	1868.4	250.42
0.016573	11.455	39.019	37.89	144.067	1373.2	580.39
0.020242	10.019	27.368	17.63	135.256	974.01	724.38
0.024724	8.6830	18.495	7.213	124.148	662.70	917.46
0.030197	7.4420	11.934	2.471	111.328	428.24	950.16
0.036883	6.2962	7.2589	0.647	97.6721	260.70	897.90
0.045049	5.2435	4.0822	0.110	83.5664	145.88	812.72
0.055023	4.2842	2.0589	0.008	69.5476	73.141	700.85
0.067206	3.4173	0.8849	0.000	56.1021	31.186	571.07
0.082085	2.6424	0.2950	−0.00	43.4943	10.302	449.83
0.100259	1.9584	0.0630	−0.00	32.1309	2.1788	326.17
0.122456	1.3647	0.0056	0.000	22.1977	0.1915	218.79
0.149569	0.8612	0.0000	0.000	13.8521	0.0005	133.14
0.182684	0.4500	−0.000	0.000	7.16729	−0.000	73.647
0.223130	0.1366	0.0000	−0.00	2.15434	0.0000	24.828
0.272532	0.1320	−0.000	0.000	2.08088	0.0000	23.734

Table 3. The perturbative coefficients entering eq. (2.5) as calculated by the MCCSM program.

B MC simulations at hadron and parton levels

The figures in this section contain data measurements and predictions for jet rates calculated in very fine binning at parton (figure 8) and hadron (figures 9, 10 and 11) levels using different MC setups. The plots with ratios of hadron level predictions to data measurements are calculated with the binning used in the measurements. The error bars in the plots are calculated from the total uncertainty of the data points. The vertical lines in the plots show the fit ranges used in this work for fits of the R_2 and R_3 observables.

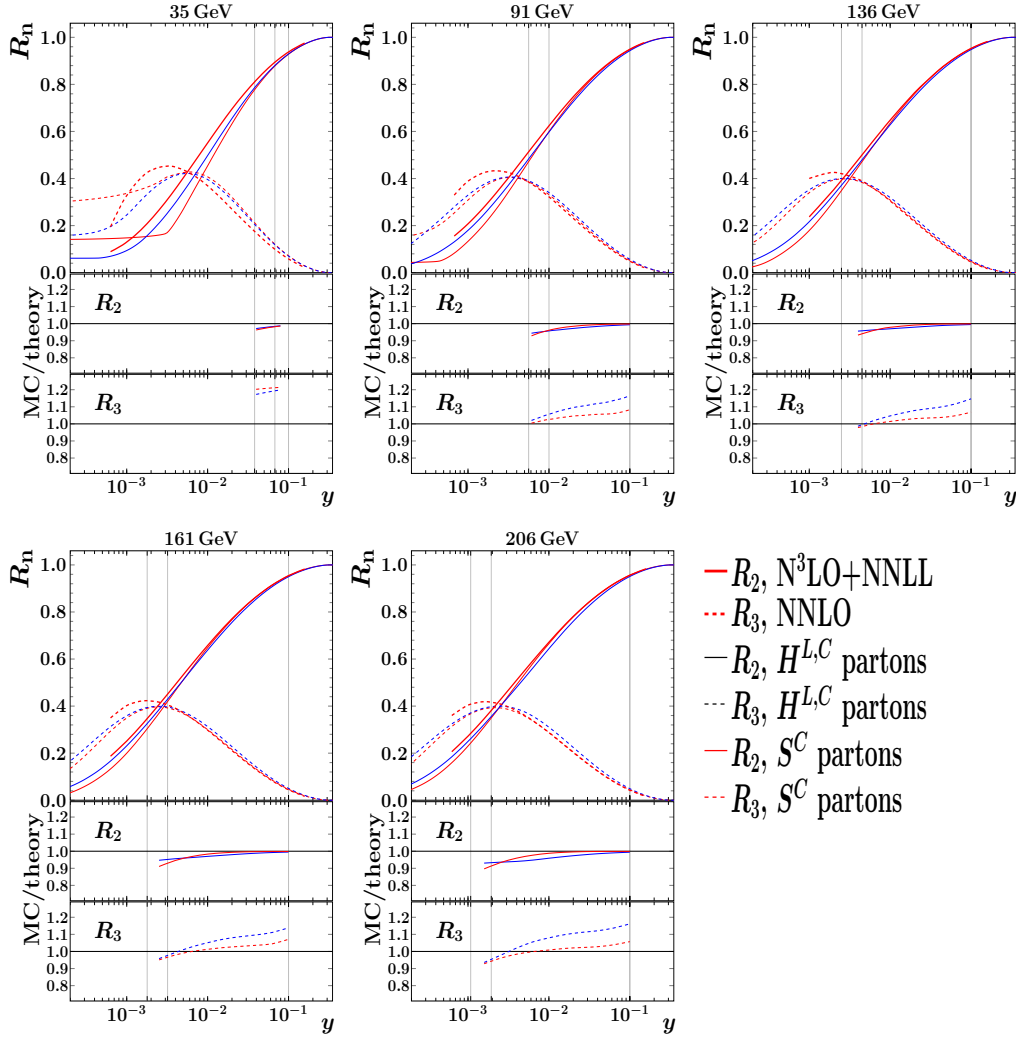


Figure 8. Selected predictions obtained with S^C , H^C and H^L MC setups at parton level and theory predictions obtained with $\alpha_S(M_Z) = 0.1181$.

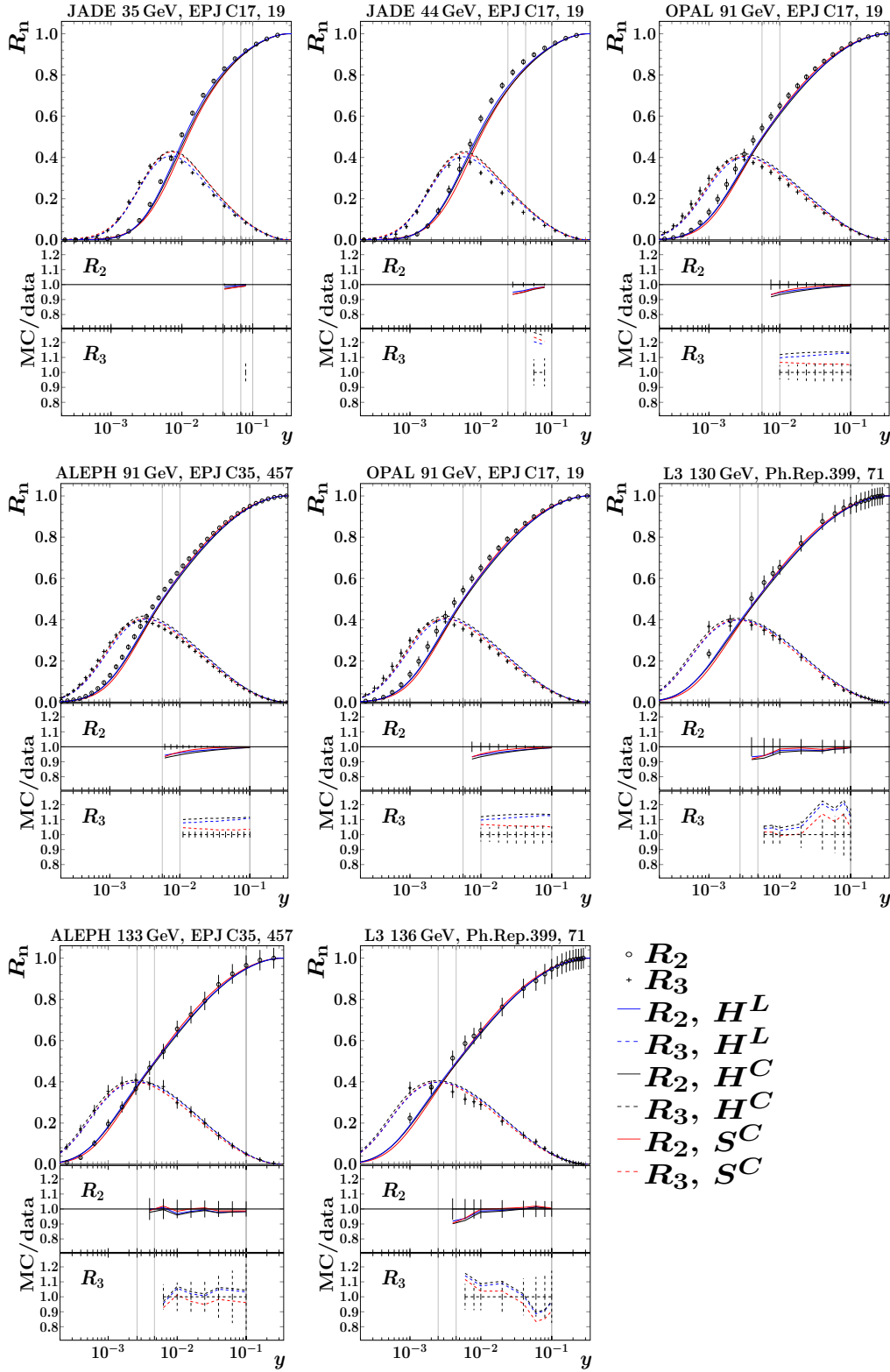


Figure 9. Predictions obtained with S^C , H^C and H^L MC setups at hadron level.

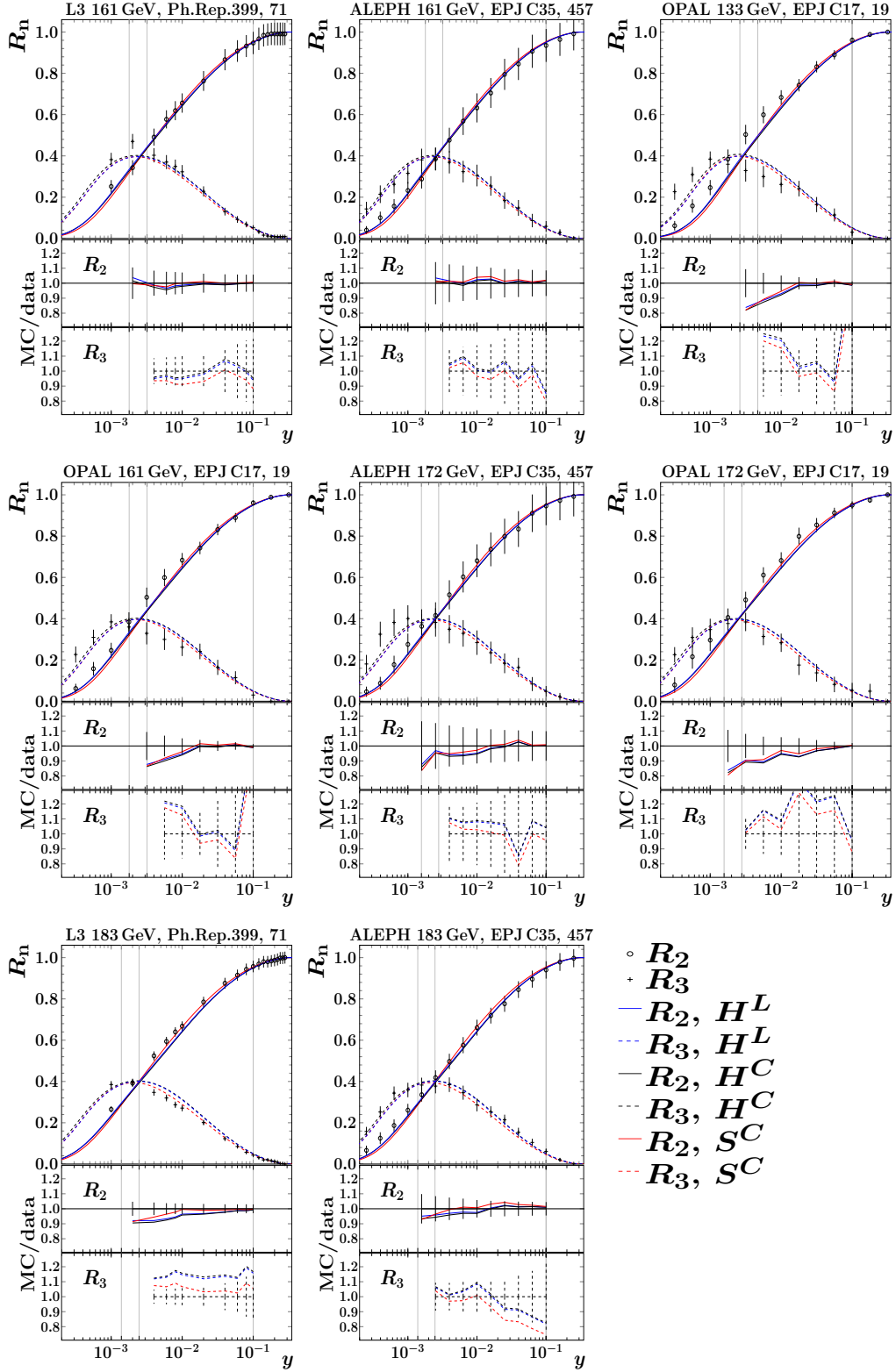


Figure 10. Predictions obtained with S^C , H^C and H^L MC setups at hadron level.

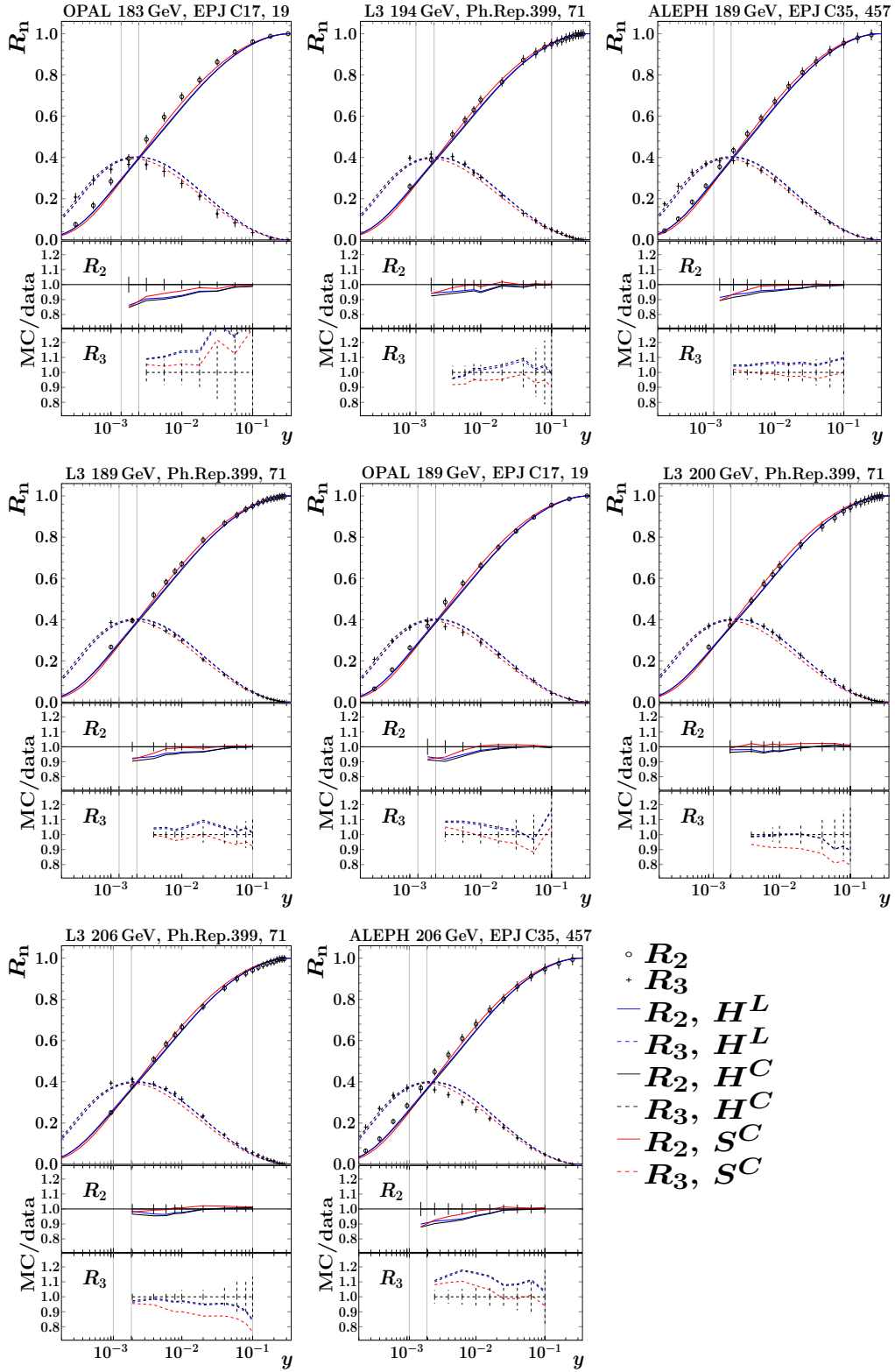


Figure 11. Predictions obtained with S^C , H^C and H^L MC setups at hadron level.

C Hadronization corrections

Figure 12 shows the $\delta\xi_1$ and $\delta\xi_2$ distributions used to model hadronization corrections of jet rates at different centre-of-mass energies. As before, the vertical lines in the plots show the fit range for reference fits of the R_2 and R_3 observables. Figure 13 shows the size of hadronization corrections obtained from the $\delta\xi_1$ and $\delta\xi_2$ values using eqs. (3.1) and (3.3).

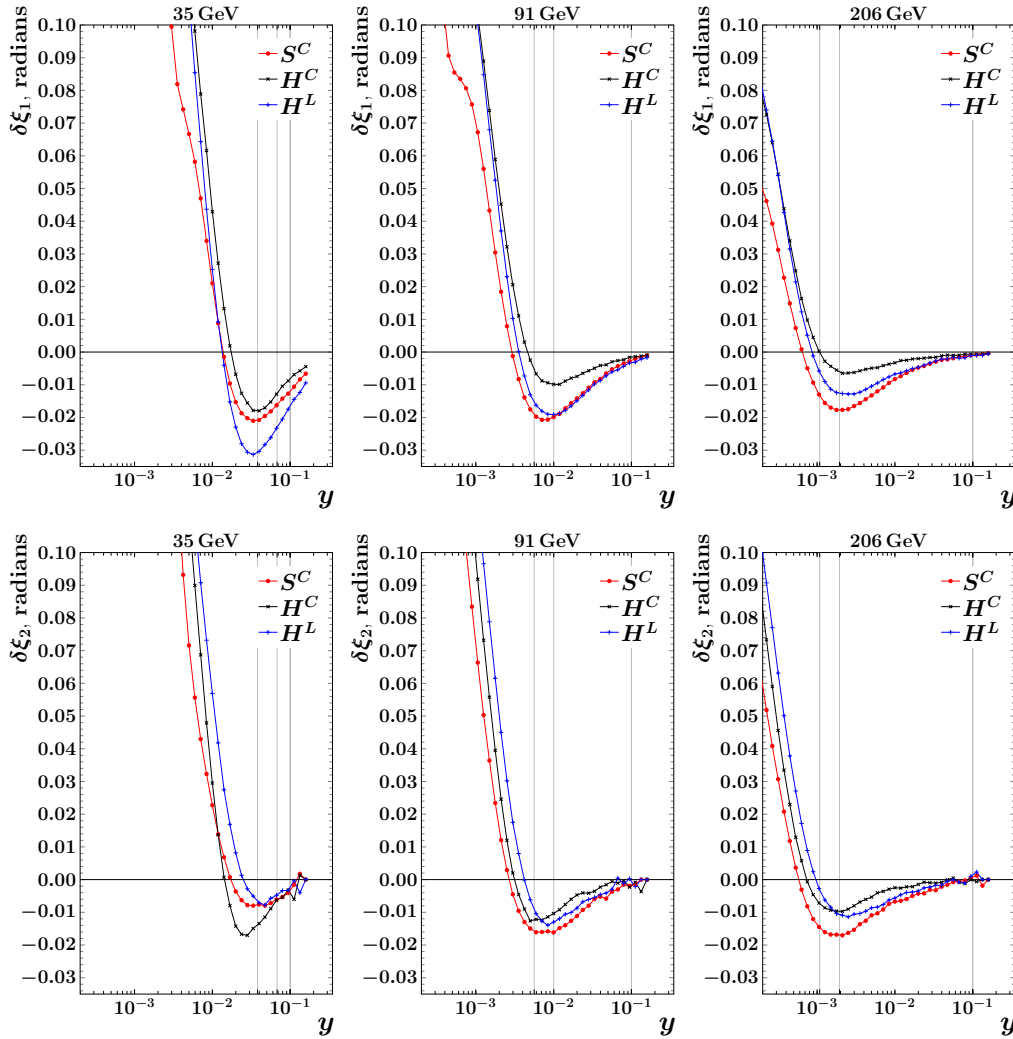


Figure 12. Hadronization corrections obtained with S^C , H^C and H^L hadronization models.

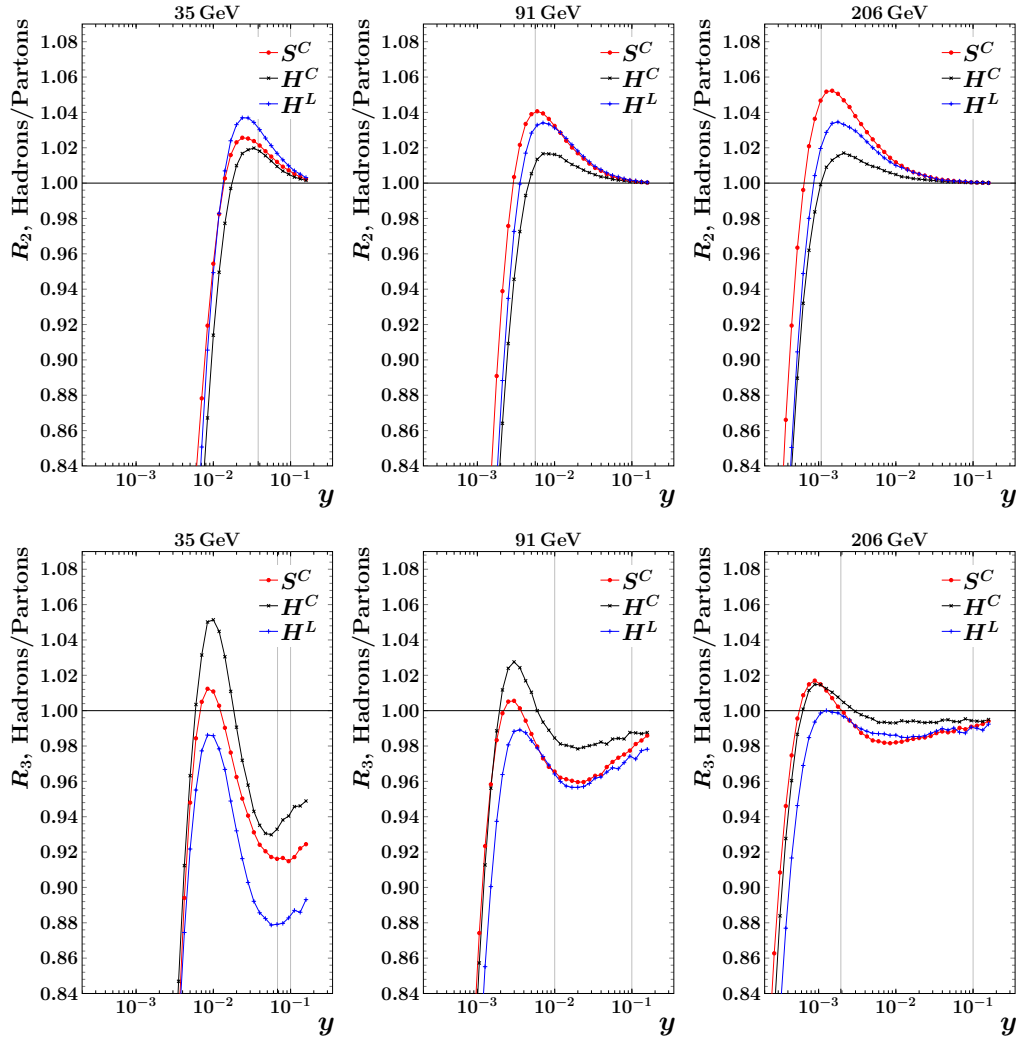


Figure 13. Hadronization corrections obtained with S^C , H^C and H^L hadronization models. For every fixed y in the fit range the deviation from unity decreases monotonically with increasing centre-of-mass energy.

D Additional fits

In this section we present results of additional fits that we have performed.

Table 4 shows a simultaneous fit of $\alpha_s(M_Z)$ from experimental data for R_2 and R_3 obtained using N³LO and N³LO+NNLL predictions for R_2 and NNLO predictions for R_3 . Three different hadronization models are used and four different choices of the fit range are shown. The reported uncertainty is only the statistical uncertainty of the fit, as given by MINUIT2. We note that these results show a significant dependence on the fit range used.

Fit ranges, log y Hadronization	N ³ LO, NNLO $\chi^2/ndof$	N ³ LO+NNLL, NNLO $\chi^2/ndof$
$[-1.75 + \mathcal{L}, -1][-1.5 + \mathcal{L}, -1]$ S^C	0.12195 ± 0.00072 120/143 = 0.84	0.12078 ± 0.00066 140/143 = 0.98
$[-2 + \mathcal{L}, -1][-1.75 + \mathcal{L}, -1]$ S^C	0.12163 ± 0.00061 153/187 = 0.82	0.12065 ± 0.00056 176/187 = 0.94
$[-2.25 + \mathcal{L}, -1][-2 + \mathcal{L}, -1]$ S^C	0.12075 ± 0.00044 208/251 = 0.83	0.11994 ± 0.00041 222/251 = 0.88
$[-2.5 + \mathcal{L}, -1][-2.25 + \mathcal{L}, -1]$ S^C	0.12143 ± 0.00043 321/331 = 0.97	0.12089 ± 0.00044 336/331 = 1.01
$[-1.75 + \mathcal{L}, -1][-1.5 + \mathcal{L}, -1]$ H^C	0.12068 ± 0.00073 126/143 = 0.88	0.11956 ± 0.00066 147/143 = 1.03
$[-2 + \mathcal{L}, -1][-1.75 + \mathcal{L}, -1]$ H^C	0.12006 ± 0.00061 163/187 = 0.87	0.11913 ± 0.00054 188/187 = 1.01
$[-2.25 + \mathcal{L}, -1][-2 + \mathcal{L}, -1]$ H^C	0.11869 ± 0.00043 221/251 = 0.88	0.11793 ± 0.00043 238/251 = 0.95
$[-2.5 + \mathcal{L}, -1][-2.25 + \mathcal{L}, -1]$ H^C	0.11845 ± 0.00045 302/331 = 0.91	0.11799 ± 0.00047 310/331 = 0.94
$[-1.75 + \mathcal{L}, -1][-1.5 + \mathcal{L}, -1]$ H^L	0.12248 ± 0.00068 121/143 = 0.85	0.12129 ± 0.00063 141/143 = 0.99
$[-2 + \mathcal{L}, -1][-1.75 + \mathcal{L}, -1]$ H^L	0.12211 ± 0.00057 155/187 = 0.83	0.12110 ± 0.00053 180/187 = 0.96
$[-2.25 + \mathcal{L}, -1][-2 + \mathcal{L}, -1]$ H^L	0.12071 ± 0.00044 209/251 = 0.83	0.11989 ± 0.00045 227/251 = 0.90
$[-2.5 + \mathcal{L}, -1][-2.25 + \mathcal{L}, -1]$ H^L	0.12041 ± 0.00044 266/331 = 0.80	0.11990 ± 0.00044 278/331 = 0.84

Table 4. Simultaneous fit of $\alpha_s(M_Z)$ from experimental data for R_2 and R_3 obtained using N³LO and N³LO+NNLL predictions for R_2 and NNLO predictions for R_3 . Three different hadronization models are used and four different choices of the fit range are shown, as given in the brackets, with $\mathcal{L} = \log_{10}(M_Z^2/Q^2)$. The fit range for R_2 is given in the first pair of brackets and the the fit range for R_3 in the second. The reported uncertainty is the fit uncertainty only, as given by MINUIT2.

Fit ranges, $\log y$ Hadronization	NNLO $\chi^2/ndof$
$[-1.5 + \mathcal{L}, -1]$ S^C	0.12176 ± 0.00113 $39/56 = 0.70$
$[-1.75 + \mathcal{L}, -1]$ S^C	0.12088 ± 0.00088 $56/86 = 0.65$
$[-2 + \mathcal{L}, -1]$ S^C	0.11996 ± 0.00074 $74/100 = 0.74$
$[-2.25 + \mathcal{L}, -1]$ S^C	0.11853 ± 0.00068 $111/150 = 0.74$
$[-1.5 + \mathcal{L}, -1]$ H^C	0.12053 ± 0.00114 $41/56 = 0.74$
$[-1.75 + \mathcal{L}, -1]$ H^C	0.11933 ± 0.00084 $61/86 = 0.70$
$[-2 + \mathcal{L}, -1]$ H^C	0.11810 ± 0.00074 $83/100 = 0.83$
$[-2.25 + \mathcal{L}, -1]$ H^C	0.11645 ± 0.00061 $125/150 = 0.83$
$[-1.5 + \mathcal{L}, -1]$ H^L	0.12257 ± 0.00112 $39/56 = 0.70$
$[-1.75 + \mathcal{L}, -1]$ H^L	0.12178 ± 0.00088 $57/86 = 0.66$
$[-2 + \mathcal{L}, -1]$ H^L	0.12076 ± 0.00068 $75/100 = 0.75$
$[-2.25 + \mathcal{L}, -1]$ H^L	0.11931 ± 0.00067 $115/150 = 0.77$

Table 5. Fit of $\alpha_s(M_Z)$ from experimental data for R_3 obtained using NNLO predictions, three different hadronization models and four different choices of the fit range, as given in the brackets, where $\mathcal{L} = \log_{10}(M_Z^2/Q^2)$. The reported uncertainty is the fit uncertainty only, as given by MINUIT2.

We also show in table 5 a fit of $\alpha_s(M_Z)$ from experimental data for R_3 only obtained using NNLO predictions, three different hadronization models and four different choices of the fit range. As before, the reported uncertainty is the fit uncertainty only, as given by MINUIT2. These results can be compared directly with results of similar analyses, such as the study of ref. [74]. Altogether we find good agreement between the fitted values reported in table 5 and the results of ref. [74].

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